

## BRIEF COMMUNICATIONS

## A CYLINDRICAL WIRE HEATED BY A CURRENT

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**ABSTRACT:** An examination is made of the radial distribution of temperature and current density in an infinitely long cylindrical wire, heated by a current allowing for the dependence of thermal and electrical conductivities on temperature.

The temperature distribution in an infinitely long cylindrical wire heated by a constant current, for constant values of the thermal conductivity  $\lambda$  and electrical conductivity  $\sigma$  has been given in [1], and for the case  $\lambda = \text{const}$ ,  $\sigma = \text{var}$  in [2].

This paper presents an approximate method of solution for the case  $\lambda = \lambda(T)$ .

The equation of energy balance for a cylindrical wire of unit length in the absence of axial overflow of heat may be written as

$$\sigma E^2 + \frac{1}{r} \frac{d}{dr} \left( r \lambda \frac{dT}{dr} \right) = 0, \quad (1)$$

where  $E$  is the longitudinal intensity of the electric field;  $r$  is the radius;  $T$  is the temperature. Using the heat conduction function

$$S = \int_0^T \lambda dT \quad [3] \quad \text{and the relative radius } \rho = (r)/R \quad (R \text{ being the wire radius}),$$

Eq. (1) may be transformed to the form

$$\sigma E^2 R^2 + \frac{1}{\rho} \frac{d}{d\rho} \left( \rho \frac{dS}{d\rho} \right) = 0. \quad (2)$$

In reducing the material function  $\alpha(S)$ , i. e., the dependence of the electrical conductivity  $\sigma$  on the thermal conductivity function  $S$ , we obtain an equation with a single parameter, considering the electric field intensity  $E$  to be constant along the wire radius. The material function  $\alpha(S)$  may be represented graphically, using the dependences  $\alpha(T)$  and  $S(T)$ . A solution is effected by linearization of  $\alpha(S)$  along the radius of the channel of the wire of the form  $\sigma = AS + B$  (the simplest case). Then the solution of (2) for  $A > 0$  becomes

$$S(\rho) = c_1 J_0(\nu) + c_2 Y_0(\nu) - B/A, \quad (3)$$

and for  $A < 0$

$$S(\rho) = c_1 I_0(\nu) + c_2 K_0(\nu) - B/A, \quad (4)$$

where  $\nu = \rho ER A^{1/2}$ ;  $J_0(\nu)$  and  $Y_0(\nu)$  are Bessel functions of zero order of the first and second kinds;  $I_0(\nu)$  and  $K_0(\nu)$  are modified Bessel function of zero order of the first and second kinds;  $c_1, c_2$ , are constants of integration.

Using the boundary conditions

$$\text{for } \rho = 1 \quad S = S_W \quad \text{and} \quad q = -\frac{1}{R} \left( \frac{dS}{d\rho} \right)_{\rho=1}, \quad (5)$$

where  $S_W$  is the value of the thermal conductivity function along the wire surface; and  $q$  is the heat flux, we obtain the radial distribution of the heat conduction function for the case  $A < 0$  (for metals):

$$S(\rho) = \frac{qR}{\nu_1 I_1(\nu_1)} I_0(\nu) - \frac{B}{A}. \quad (6)$$

Here  $\nu_1 = ER(A)^{1/2}$  is determined from solution of the transcendental equation

$$\frac{\nu_1 I_1(\nu_1)}{I_0(\nu_1)} = \frac{qR}{S_{CT} + B/A}, \quad (7)$$

where the quantities  $q, R, A, B$  and  $S_W$  are assumed to be given, and  $I_1(\nu_1)$  is a modified Bessel function of the first order and the first kind.

Thereafter we may determine the value  $E$ , and knowing the heat flux  $q$  and  $E$ , find the current. The radial distribution of current density (the density of the internal source of heat generation) may be expressed as

$$j(\rho) = \frac{\sigma_W E}{I_0(\nu_1)} I_0(\nu), \quad (8)$$

where  $\sigma_W$  is the electrical conductivity at the temperature of the wire surface.

Expression (8) illustrates the phenomenon of current-density "thermal skin effect" due to decrease of the electrical conductivity with increase of temperature. The radial distribution of temperature may be obtained from the profile  $S(\rho)$ , using the graphical dependence of  $S(T)$ .

If we use the boundary conditions

$$\begin{aligned} \text{for } \rho = 1 \quad S &= S_W, \\ \text{for } \rho = 0 \quad S &= S_0, \end{aligned} \quad (9)$$

we obtain the following expressions for the distribution of the heat conduction function and of the specific heat flux through the wire surface,

$$S(\rho) = (S_0 + B/A) I_0(\nu) - B/A, \quad (10)$$

$$q = \sigma_0 \nu_1 I_1(\nu_1)/RA, \quad (11)$$

where  $\sigma_0$  is the electrical conductivity at the temperature on the axis, and  $\nu_1 = ER(A)^{1/2}$  is determined from (10) by substitution of the boundary conditions.

The accuracy of the approximate solution will depend on the nature of the approximation  $\alpha(S)$ . For increased accuracy of solution the dependence  $\alpha(S)$  may be approximated by several straight lines [4].

It follows from the foregoing examination that for given temperatures on the axis and on the wire surface the required electrical power  $E$  does not depend on the wire radius. Then we will find constant temperature profiles with respect to the relative radius. The condition of similarity in this case is written in a similar way to the condition for a positive arc column [4]:  $ER = \text{const}$  for  $EI = \text{const}$  (i. e., with constant boundary conditions), where  $I$  is the current.

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